

# Micro-local analysis of hyperbolic dynamics. Part I: Anosov flows

F. Faure (Grenoble) with M. Tsujii (Kyushu),

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(this notes, based on [arXiv:1706.09307](https://arxiv.org/abs/1706.09307) can be downloaded on the [web page of  
frédéric faure](#))

Part II will be “Geodesic (or contact) Anosov flows”.

# Outline

- ① **Anosov vector field**  $X : C^\infty(M) \rightarrow C^\infty(M)$ ,  
**transfer operator (pull-back)**  $e^{tX} u = u \circ \phi^t$  and examples of  
**Ruelle-Pollicott discrete spectra.**
- ② Micro-local analysis with wave-packets transform  $\mathcal{T}$  (an isometry, i.e.  
 $\mathcal{T}^* \mathcal{T} = \text{Id}$ ):

$$\begin{array}{ccc} L^2(M) & \xrightarrow{e^{tX}} & L^2(M) \\ \tau \downarrow & & \tau \downarrow \\ L^2(T^*M) & \xrightarrow{\mathcal{T} e^{tX} \mathcal{T}^*} & L^2(T^*M) \end{array} \quad (1)$$

- ① **Trace formula** for  $A \subset T^*M$ :  $\text{Tr}(\mathcal{T}^* \chi_A \mathcal{T}) \asymp \frac{1}{(2\pi)^{\dim M}} \text{Vol}(A)$ . :“**uncertainty principle**”, “density of information is finite”.
- ② **Theorem of propagation of singularities**: The Schwartz kernel of  $\mathcal{T} e^{tX} \mathcal{T}^*$  decays very fast outside the graph of  $(D\phi^t)^* : T^*M \rightarrow T^*M$ .
- ③ **Anisotropic Sobolev space**  $\mathcal{H}_W(M)$ : for  $u \in C^\infty(M)$ , define

$$\|u\|_{\mathcal{H}_W(M)} := \|W \mathcal{T} u\|_{L^2(T^*M)}$$

with a **weight function**  $W : T^*M \rightarrow \mathbb{R}^+$ , that is (Lyapunov-escape function for  $(D\phi^t)^*$ )

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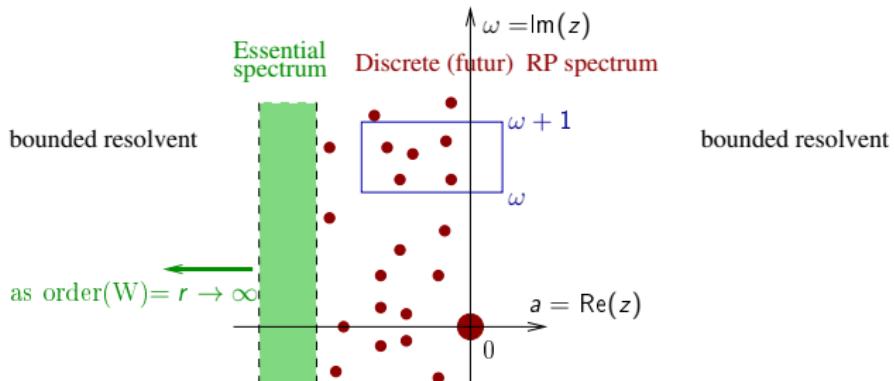
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## ③ Results

### ① Group properties of $e^{tX} : \mathcal{H}_W(M)$ ○

and  $X$  has “futur/past Ruelle-Pollicott **discrete spectrum**”:



Rem: eigenvalue  $z_j = a_j + i\omega_j \Rightarrow e^{tz_j} = \underbrace{e^{ta_j}}_{\text{decay if } a_j < 0, t \rightarrow +\infty, \text{ oscillation}} \underbrace{e^{it\omega_j}}_{\cdot}$ .

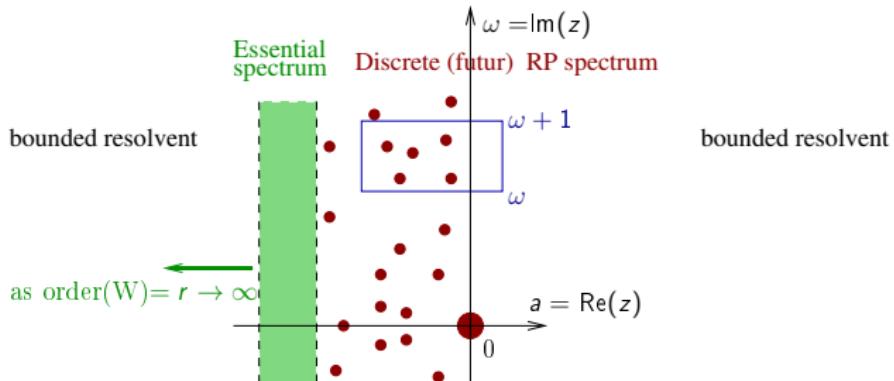
- ② “**Fractal Weyl law**”: upper bound for the density of eigenvalues  $\leq Cw^{\frac{\dim M - 1}{1 + \beta_0}}$ , with  $\beta_0$ : Hölder exp. of  $E_u \oplus E_s$ .
- ③ “**Wave front set**”: where  $(\mathcal{T}u)(\rho), \rho \in T^*M$  is non negligible for an eigenfunction  $u$  of  $X$ .
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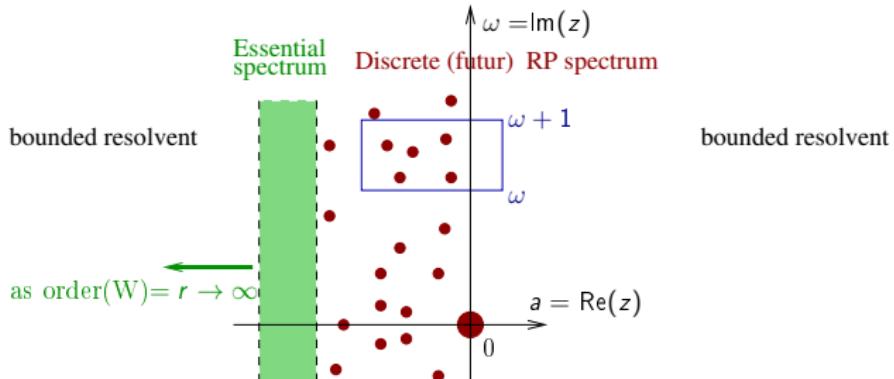
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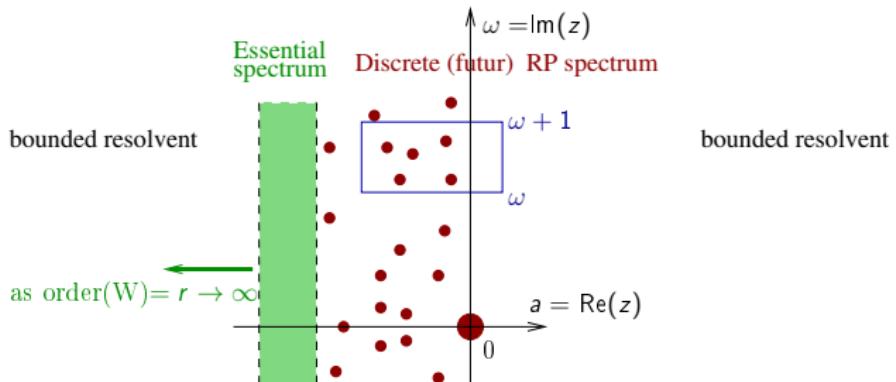
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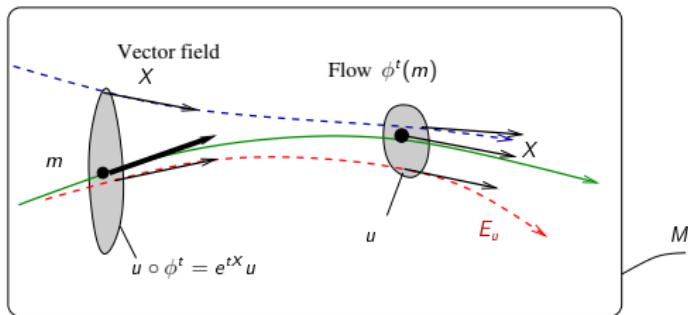
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- Series of work and interesting recent activity around Ruelle-Pollicott resonances in hyperbolic dynamics: Ruelle, Bowen 70', Pollicott 86, Rugh 90', Blank, Keller, Liverani 2002, Gouëzel, Liverani 2005, Naud 2005, Baladi, Tsujii 2005,2008, Butterley-Liverani 2007, Roy-Sjöstrand-F. 2008, Datchev-Dyatlov-Zworski, Dyatlov-Guillarmou 2014 for Axiom A flows. Dang-Rivi  re 2016 for Morse-Smale flows, Bonthonneau-Weich 2017 for cusps, Guillarmou-Hilgert-Weich 2018 for symmetric spaces, Baladi-Demers-Liverani 2018 for Sina   billiards.

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# PART 1: Deterministic dynamics: vector field $X$ and flow $\phi^t$



On a closed manifold  $M$ , let  $X \in C^\infty(M; TM)$  be a **smooth vector field** i.e. first order diff. operator

$$X \equiv \sum_{j=1}^{\dim M} X_j(y) \frac{\partial}{\partial y_j}$$

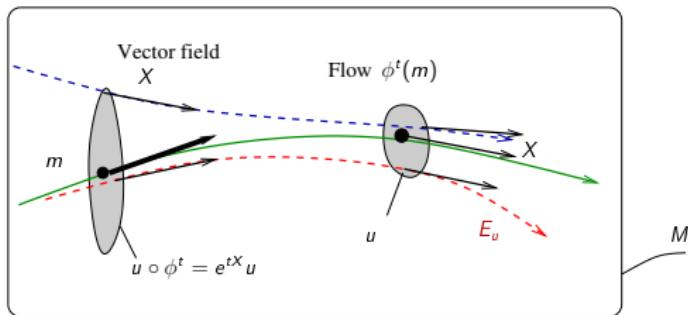
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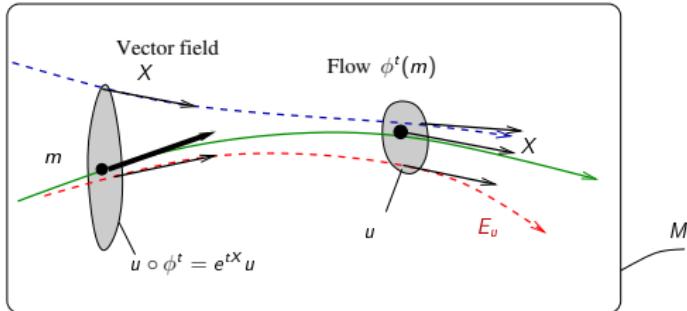
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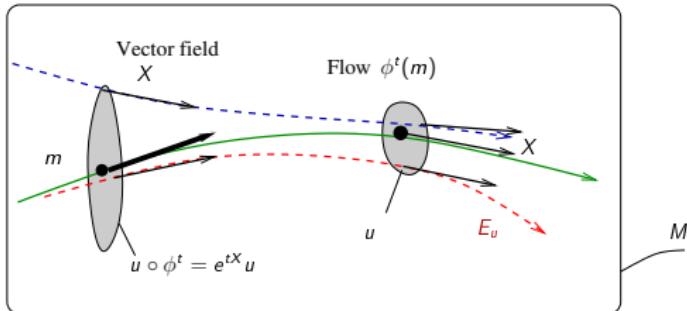
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- pushes forward probability distributions because

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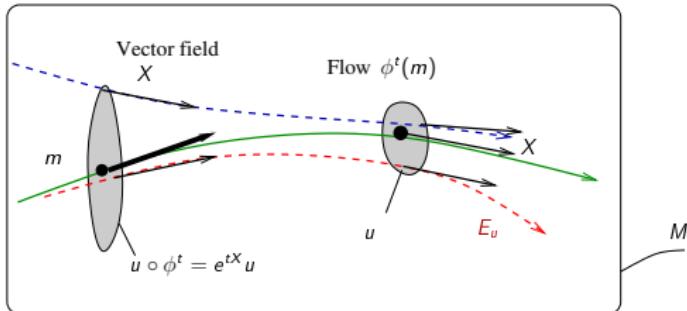
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## More general evolution of sections of $F \rightarrow M$

We may consider  $F \rightarrow M$  a **vector bundle** and  $X_F$  a first order diff. operator on sections

$$X_F : C^\infty(M; F) \rightarrow C^\infty(M; F)$$

s.t. **Leibniz condition** holds (i.e.  $X_F$  is an “extension” of  $X$ ):

$$\forall f \in C^\infty(M), u \in C^\infty(M; F), \quad X_F(fu) = X(f)u + fX_F(u).$$

- Example: tensor bundle  $F = TM \otimes \dots \otimes T^*M$ .  $X_F$  is the Lie derivative and  $e^{tX_F}$  transports tensor fields.
- Example: on the trivial bundle  $F = M \times \mathbb{C}$ ,  $X_F = X + V$  with  $V \in C^\infty(M)$ : “**Gibbs potential**”, then

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# Anosov flow (or uniformly hyperbolic flow)

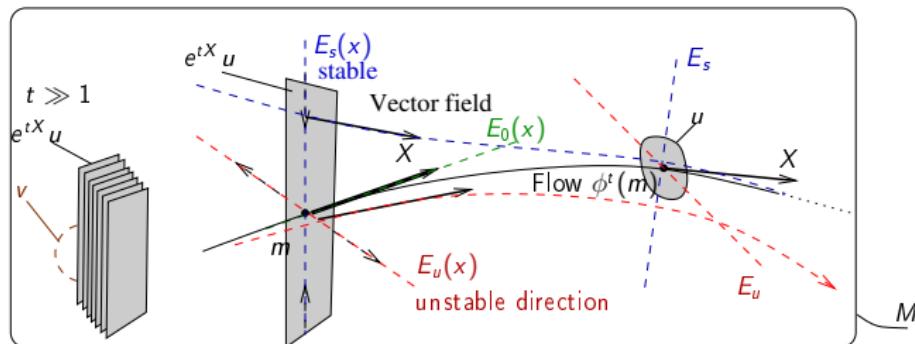
## Definition

Vector field  $X$  is **Anosov** if

$$\forall m \in M, \quad T_m M = E_u(m) \oplus E_s(m) \oplus \underbrace{E_0(m)}_{\mathbb{R}X}$$

is  $D\phi^t$ -invariant, continuous, s.t.  $\exists g, \exists C > 0, \lambda > 0, \forall t \geq 0, m \in M,$

$$\left\| D\phi^t_{/E_s(m)} \right\|_g \leq Ce^{-\lambda t}, \quad \left\| D\phi^{-t}_{/E_u(m)} \right\|_g \leq Ce^{-\lambda t},$$

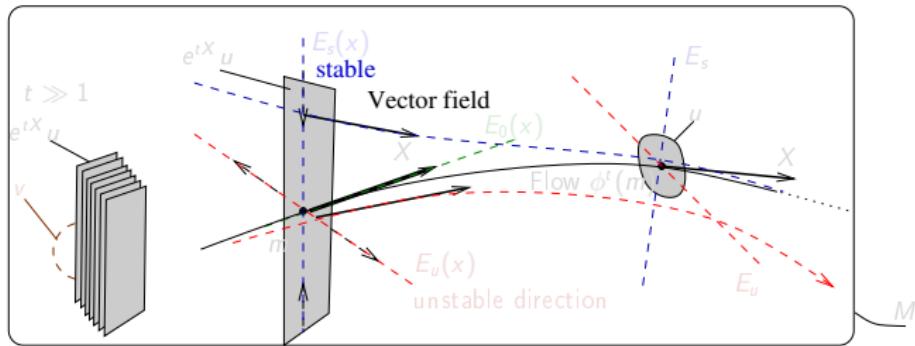


## Remarks

- Anosov property is **stable** under any (small  $C^1$ ) perturbation of  $X$ .
- maps  $m \rightarrow E_u(m), E_s(m), E_u(m) \oplus E_s(m)$  are Hölder-continues with some respective exponents  $0 < \beta_u, \beta_s, \beta_0 \leq 1$ .

## Question

Description of long time behavior of the dynamics  $\langle v | e^{tX} u \rangle$ ? i.e. **discrete spectrum of  $X$  (or  $X_F$ )?** i.e. See *movie2, movie3*.

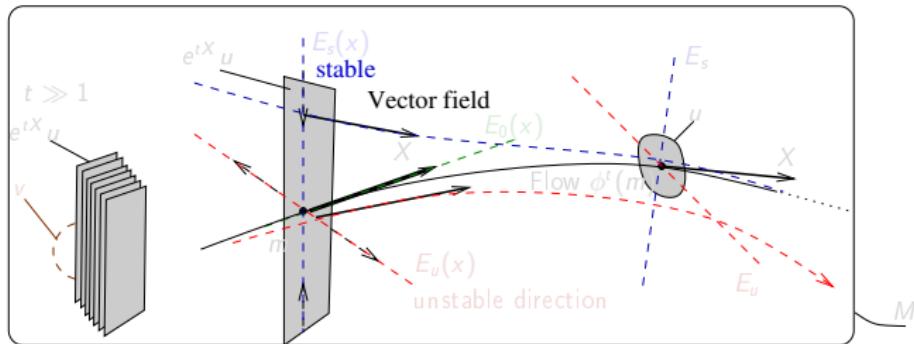


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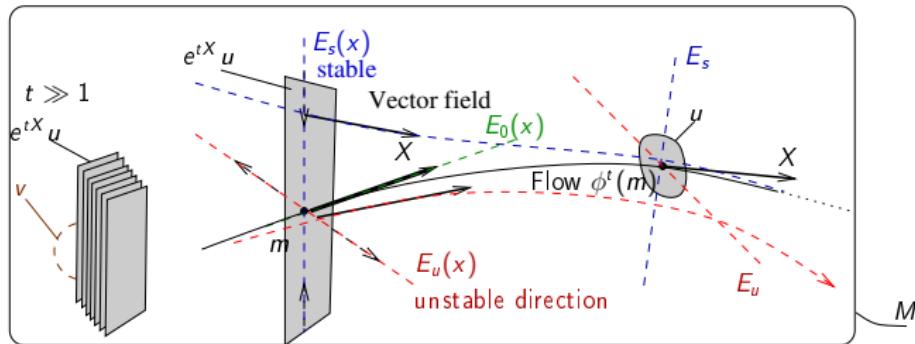


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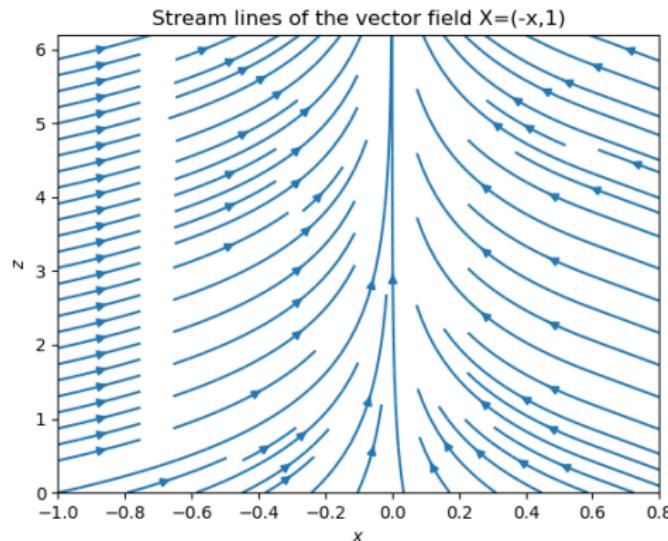


Example 1: hyperbolic toy model. (proofs later or [here](#)).

$M = \mathbb{R}_x \times (\mathbb{R}_z / (2\pi\mathbb{Z}))$ : cylinder

$$X = -x\partial_x + \partial_z. \quad \text{Flow: } \phi^t(x, z) = (e^{-t}x, z + t).$$

Trapped set or non wandering set: a single periodic orbit at  $x = 0$ :

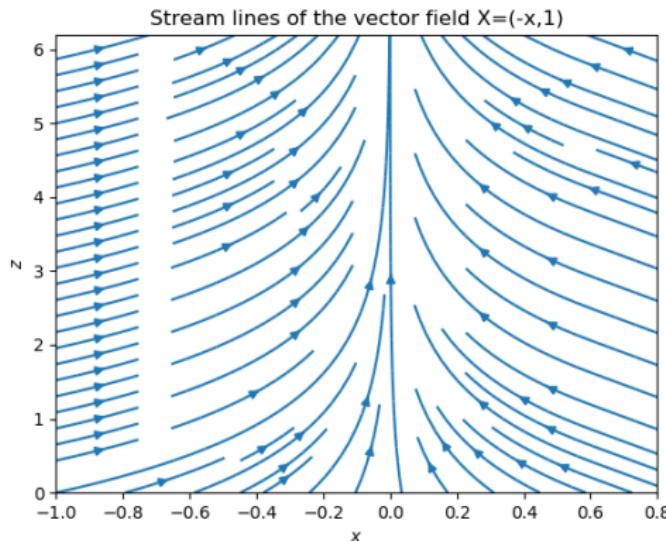


Example 1: hyperbolic toy model. (proofs later or [here](#)).

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- Let the Lyapunov (escape) function

$$W(x, \xi) = \frac{\langle \sqrt{h}\xi \rangle^r}{\langle \sqrt{h}x \rangle^r}$$

with  $0 < h \ll 1$ ,  $r \geq 0$ , (Notation:  $\langle x \rangle := |x|$  if  $|x| \geq 1$ , otherwise  $\langle x \rangle = 1$ .) and let

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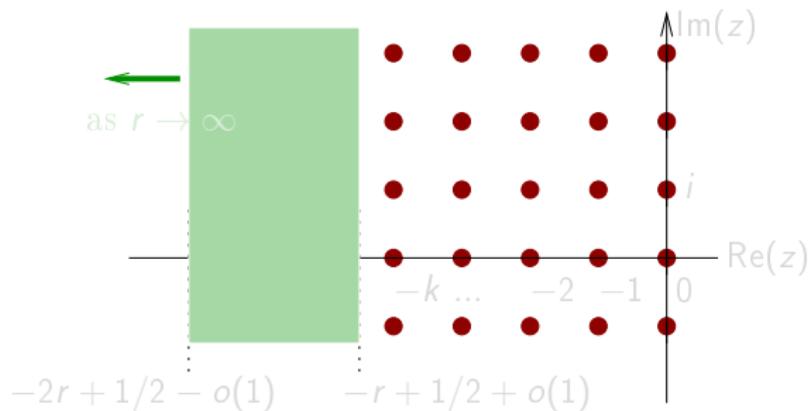
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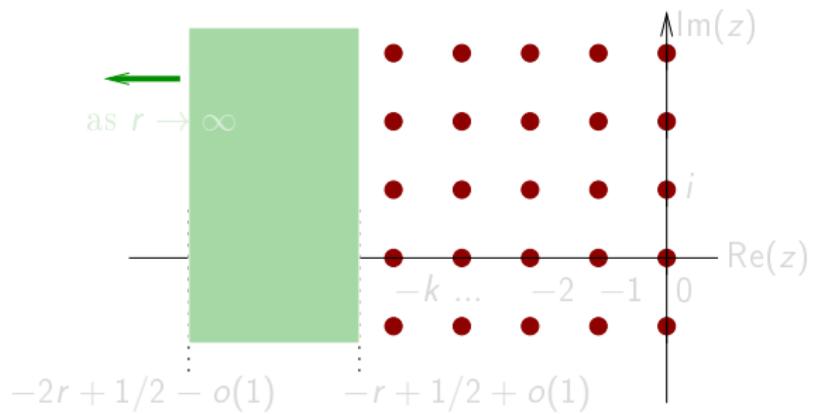
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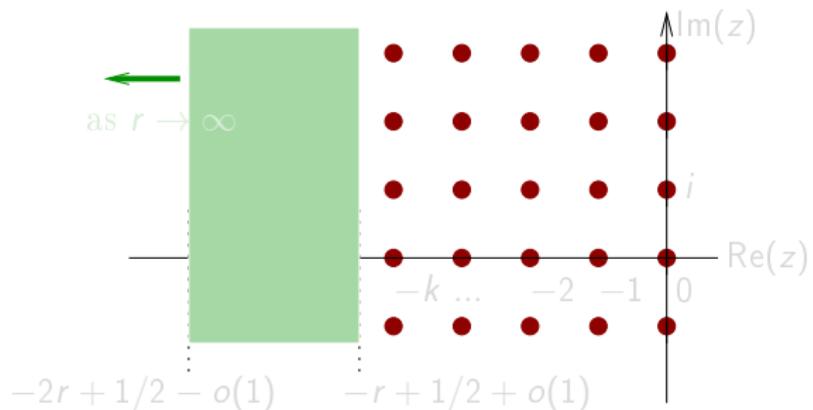
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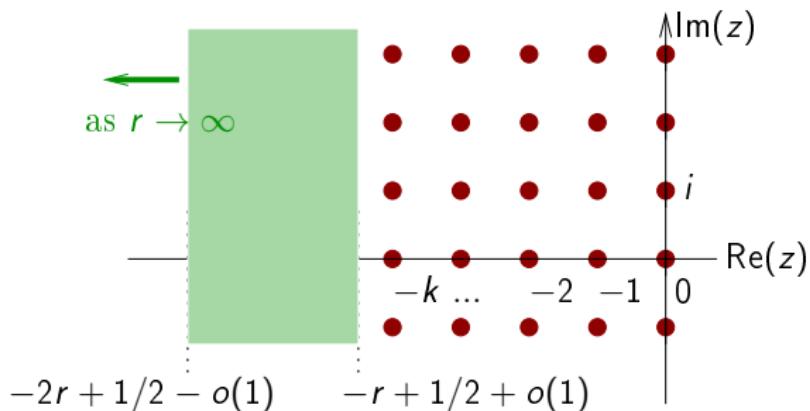
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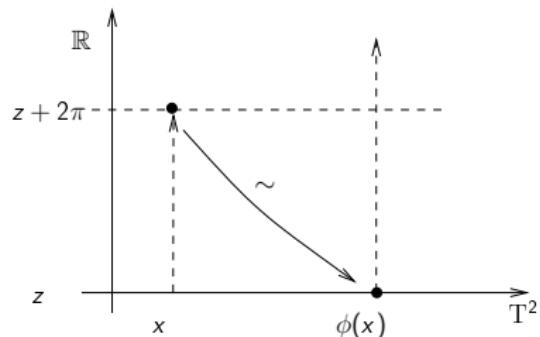
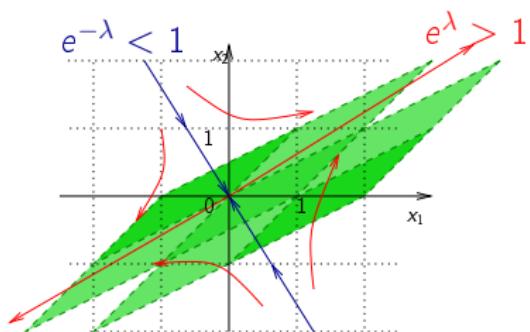
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## Example 2: suspension of a cat map

Let  $M := (\mathbb{T}^2 \times \mathbb{R}) / \sim$  with  $(x, z + 2\pi) \sim (\phi(x), z)$  with

$$\phi = \begin{pmatrix} 2 & 1 \\ 1 & 1 \end{pmatrix} : \mathbb{T}^2 \rightarrow \mathbb{T}^2 \quad : \text{cat map}$$

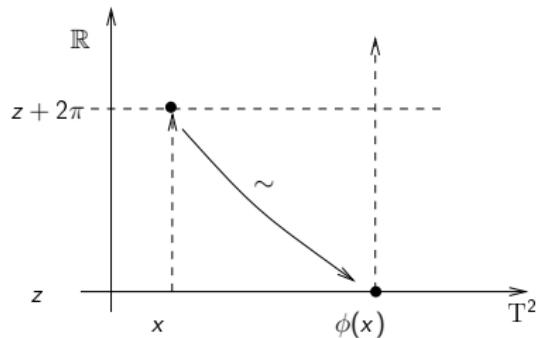
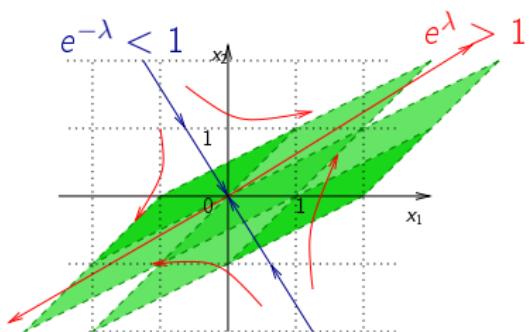


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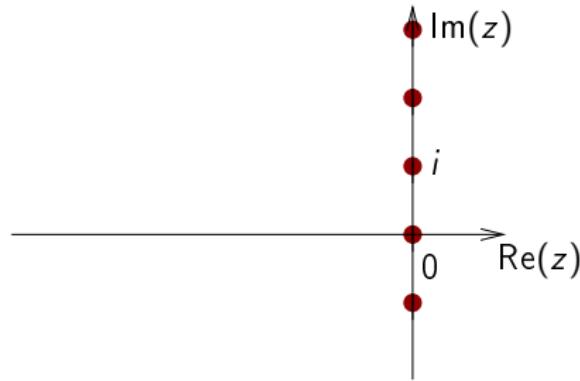
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Because for Fourier modes  $\varphi_k(x) = e^{i2\pi kx}$ ,  $\langle \varphi_{k'} | \varphi_k \circ \phi^t \rangle = 0$  for  $t \gg 1$ , except for  $k' = k = 0$ .

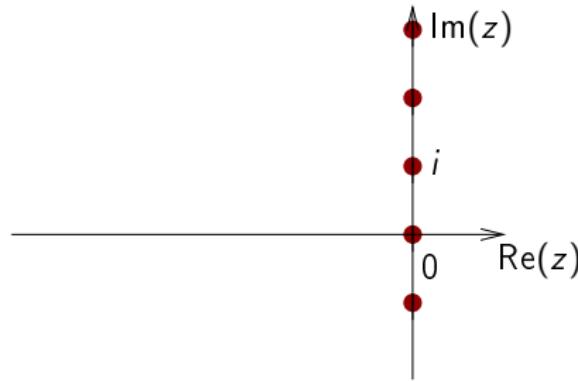
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## Example 3. Geodesic flow on hyperbolic surface (still special but less trivial)

(ref: Flaminio-Forni 02, Dyatlov-F-Guillarmou 14 for  $\Gamma \backslash SO_{1,n} / SO_{n-1}$ ).

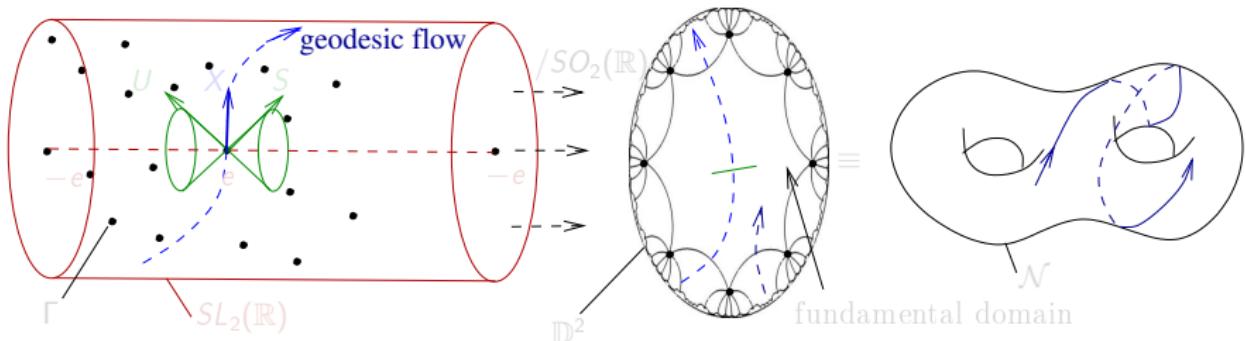
$M = \Gamma \backslash SL_2(\mathbb{R})$  smooth compact, with  $\Gamma$  cocompact subgroup.

$sl_2(\mathbb{R})$  algebra:

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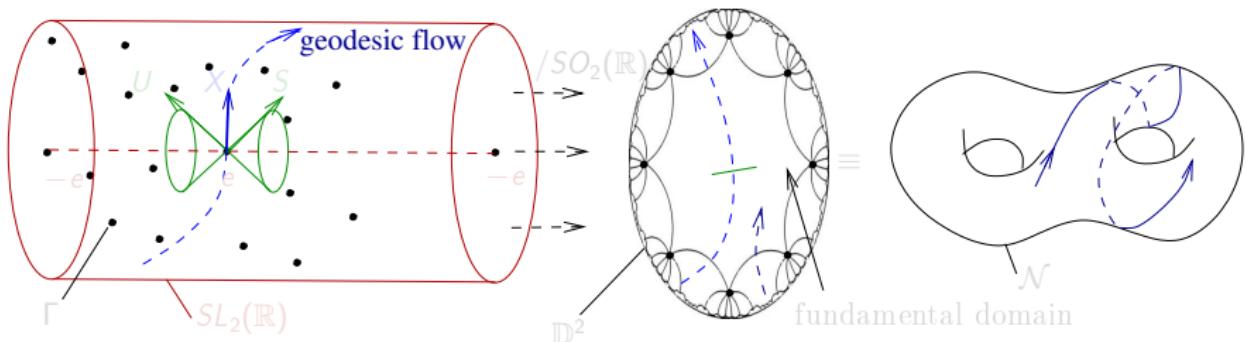
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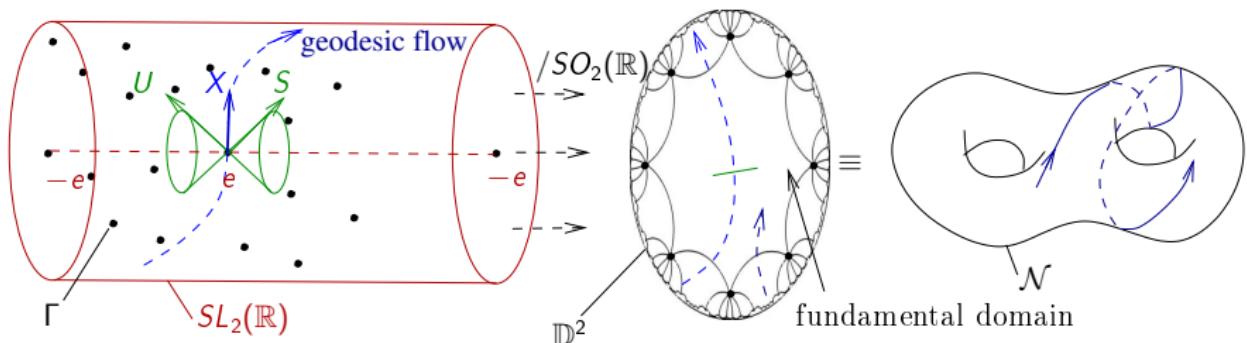
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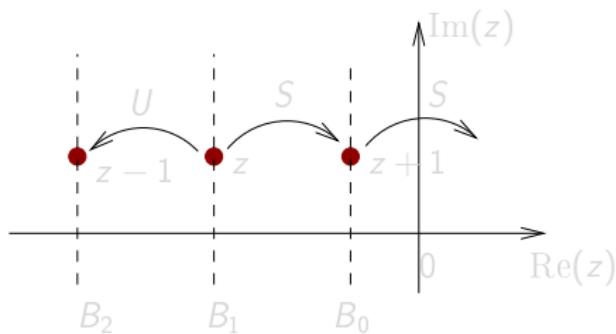
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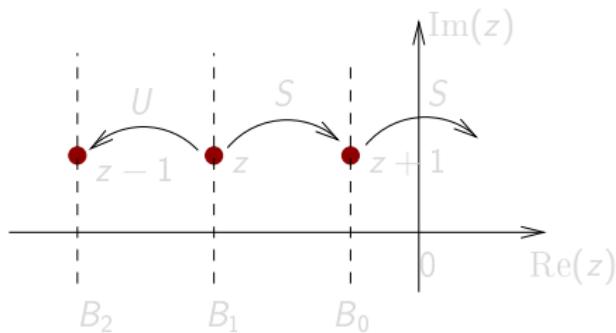
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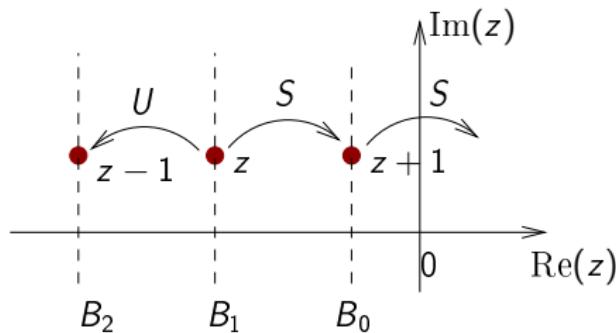
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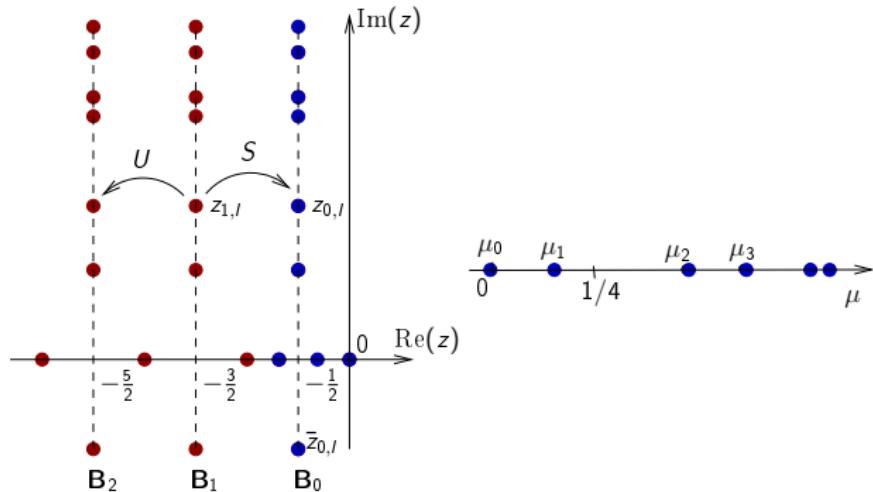
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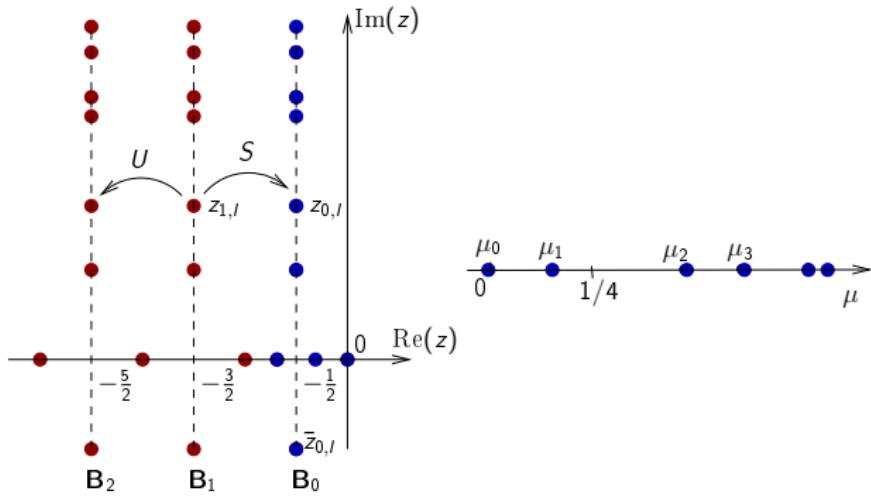


- Rem: if  $\mu_1 < \frac{1}{4}$ , the exponential rate for mixing is  $e^{(-\frac{1}{2} + \sqrt{\frac{1}{4} - \mu_1})t}$ .
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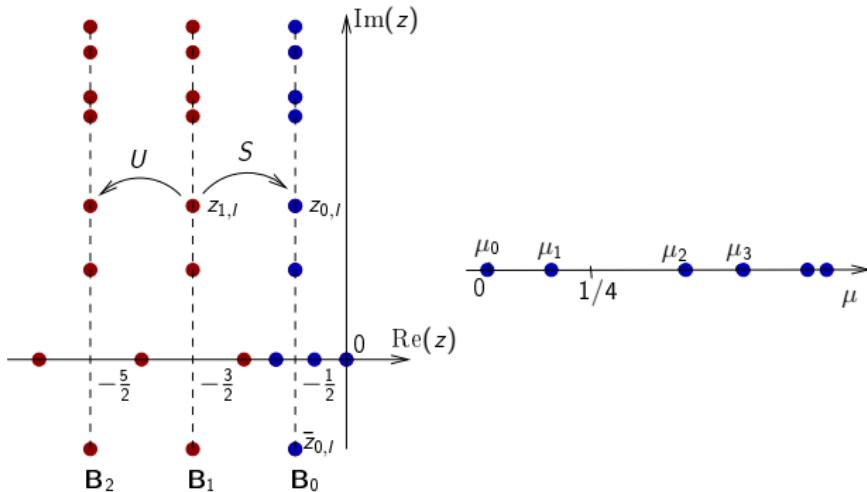


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## Example 4, from numerical computation.

Consider the partially expanding map on  $(\mathbb{R}/\mathbb{Z})_x \times \mathbb{R}_z$

$$\phi : \begin{cases} x & \rightarrow 2x \mod 1 \\ z & \rightarrow z + \sin(2\pi x) \end{cases}$$

If  $u(x, z) = v(x) e^{i\omega z}$  with  $\omega \in \mathbb{R}$ , then

$$(u \circ \phi)(x, z) = (\mathcal{L}_\omega v)(x) e^{i\omega z}$$

with

$$(\mathcal{L}_\omega v)(x) = e^{i\omega \sin(2\pi x)} v(2x).$$

Discrete RP spectrum in  $H^{-r}(S^1)$ ,  $r \gg 1$ :

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The only obvious is  $z_0(0) = 0$ . See **movie** of  $(z_j(\omega))_j$ , **movie** of  $e^{z_j(\omega)}$ , for  $\omega \in \mathbb{R}$ .

## Example 4, from numerical computation.

Consider the partially expanding map on  $(\mathbb{R}/\mathbb{Z})_x \times \mathbb{R}_z$

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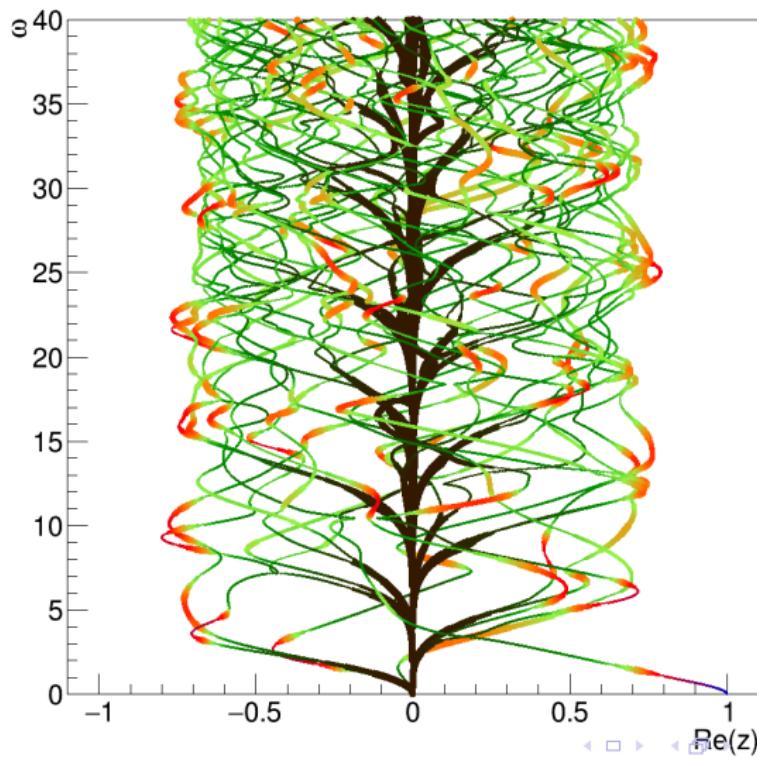
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## Tree of Ruelle-Pollicott resonances

For example 4, here is the spectrum of  $\mathcal{L}_\omega$ ,  $\operatorname{Re}(e^{z_j(\omega)})$  for  $j \in \mathbb{N}$  and  $\omega \in [0, 40]$ .  
(colors are related to  $|e^{z_j(\omega)}|$ )



## PART 2: Microlocal analysis of $X : C^\infty(M) \rightarrow C^\infty(M)$

- Objective: **understand the discrete spectrum of an Anosov vector field**  
$$X = -\sum_{j=1}^{\dim M} X_j(x) \frac{\partial}{\partial x^j}$$
 on  $M$
- We will use **wave-packet transform**, quantization, also called “FBI, wavelet, Bargmann, Anti-Wick, Wick, Toeplitz, Coherent-states” quantization. Wave-packet calculus is equivalent to the usual Weyl quantization and PDO calculus but (more) convenient for Hölder regularity.
- We will observe “**quantum scattering on a compact trapped set  $E_0^*$** ” in  $T^*M$ . From Helffer-Sjöstrand like analysis (86), we will obtain a **discrete spectrum of “Ruelle resonances” in suitable anisotropic Sobolev spaces**.

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# Wave packets

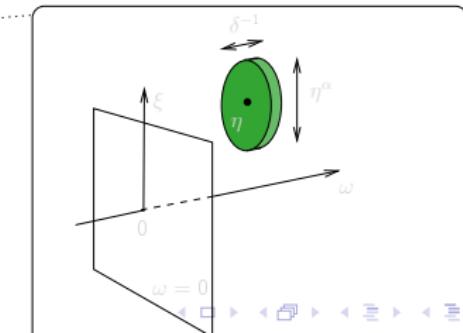
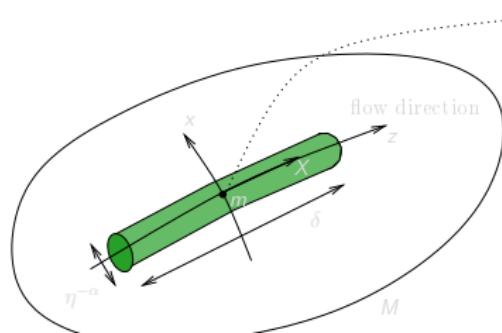
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# Wave packets

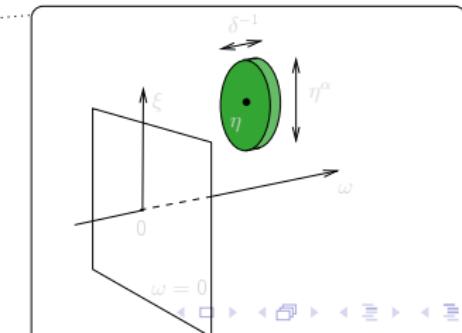
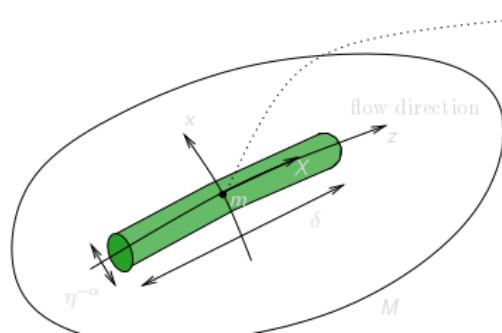
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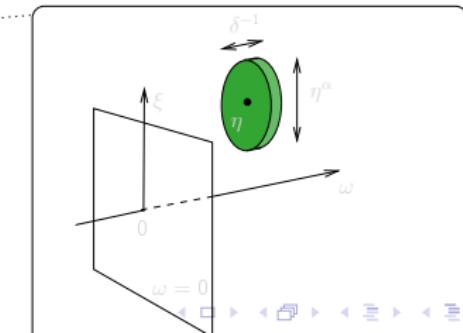
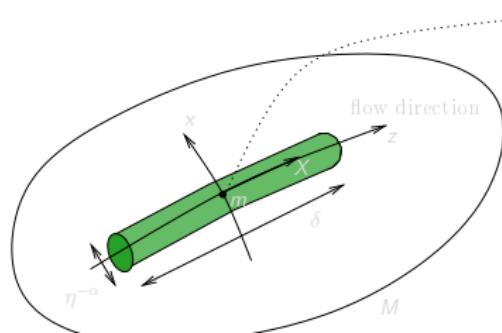
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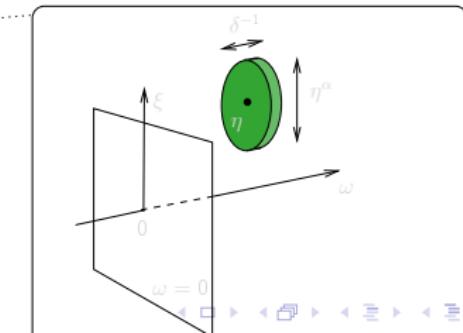
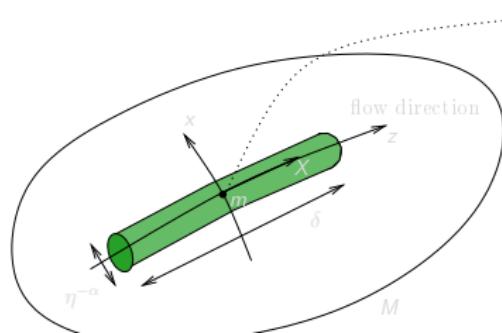
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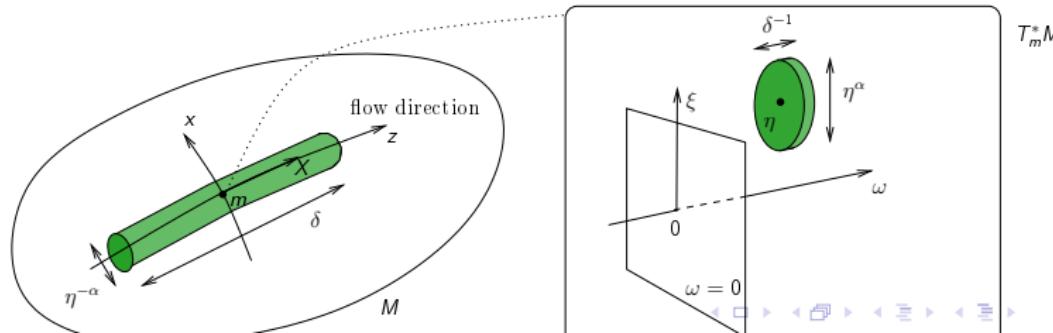
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# Wave packets and phase space metric $g$

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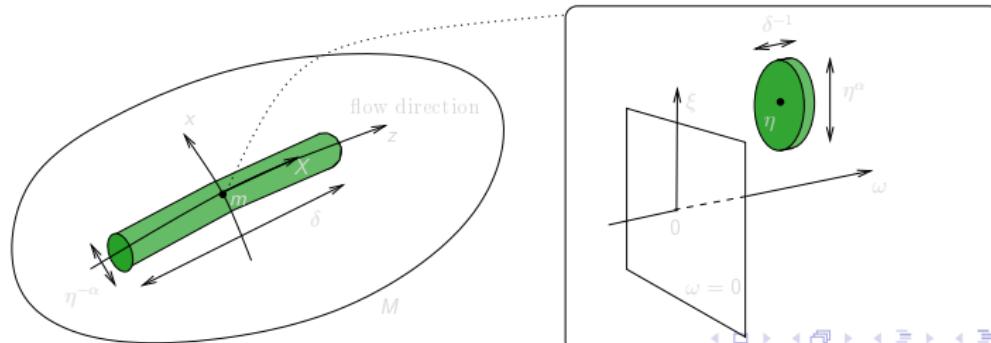
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$\forall N, \exists C_N > 0,$

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with the **metric  $g$**  on  $T^*M$ , compatible with  $\Omega = dy \wedge d\eta$ :

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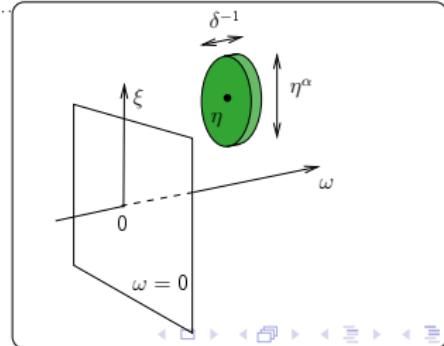
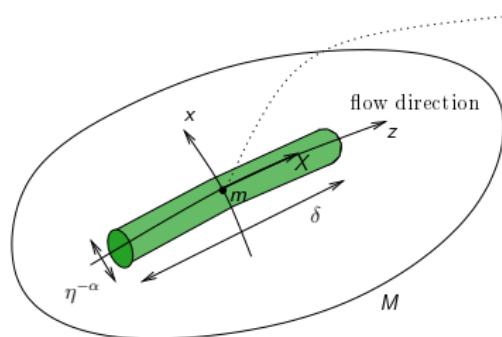
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# Phase space metric $g$

## Lemma (for microlocal analysis on manifolds)

If  $\varphi : M \rightarrow M$  is a **local “flow-diffeomorphism”** i.e.

$$\varphi : x' = f(x), z' = z + g(x) \Leftrightarrow X = \partial_z = \partial_{z'}$$

and  $\tilde{\varphi} : T^*M \rightarrow T^*M$  is the induced map, then

$$\tilde{\varphi}^* g \underset{\text{unift}/\rho}{\asymp} g \Leftrightarrow \alpha \geq \frac{1}{2}.$$

Rem: Sasaki metric on  $T^*M$ , has  $\alpha = 0$ , is not uniformly invariant by the flow.

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Proof: Want to show  $\exists C > 0, \forall \rho \in T^*M, \|D\tilde{\varphi}\|_g \leq C$ .

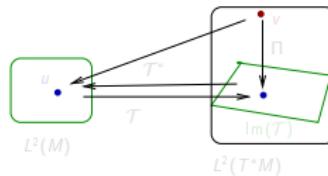
Consider  $\tilde{f}(x, \xi) = (f^{-1}(x), (Df_x)^* \xi)$ .

$$D\tilde{f} \equiv \begin{pmatrix} Df_x^{-1} & 0 \\ D(Df_x)^* \xi & (Df_x)^* \end{pmatrix}$$

$$\|D\tilde{f}\|_g \leq C + C(\langle \xi \rangle^\alpha)^{-1} \langle \xi \rangle^{-\alpha} \leq C \text{ if } \alpha \geq \frac{1}{2}.$$

**Wave packet transform** (or FBI, wavelet, Bargmann, Anti-Wick... transform)  
(Abuse of notations that forget charts and partitions of unity.)

$$\mathcal{T} : \begin{cases} C^\infty(M) & \rightarrow \mathcal{S}(T^*M) \\ u(y') & \rightarrow (\mathcal{T}u)(\rho) := \langle \varphi_\rho | u \rangle_{L^2(M)} \end{cases}$$

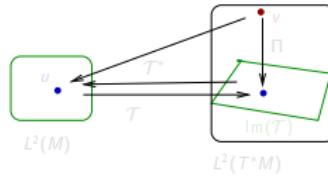


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$$\mathcal{T}^* \circ \mathcal{T} = \text{Id}_{C^\infty(M)}$$



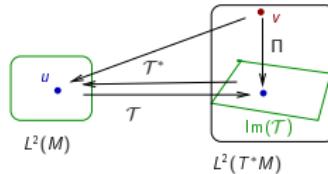
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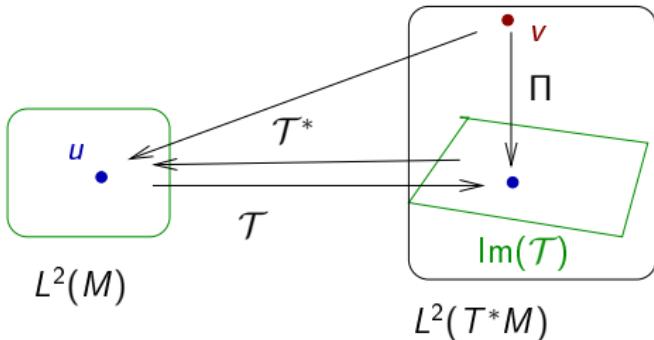
with the  $L^2(M) \rightarrow L^2(T^*M; \frac{1}{(2\pi)^{\dim M}} d\rho)$ -adjoint is

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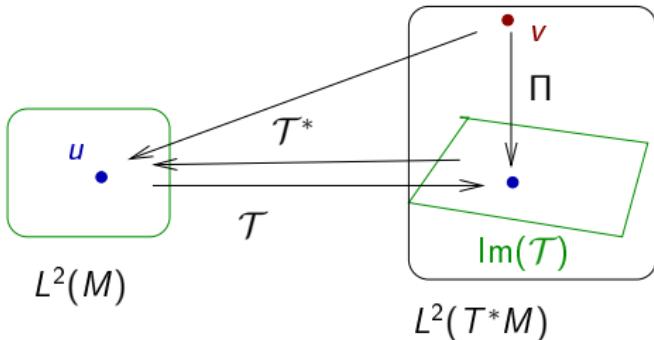
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- $\Pi = \mathcal{T} \circ \mathcal{T}^* : L^2(T^*M) \rightarrow \text{Im}(\mathcal{T})$  is an **orthogonal projector**.

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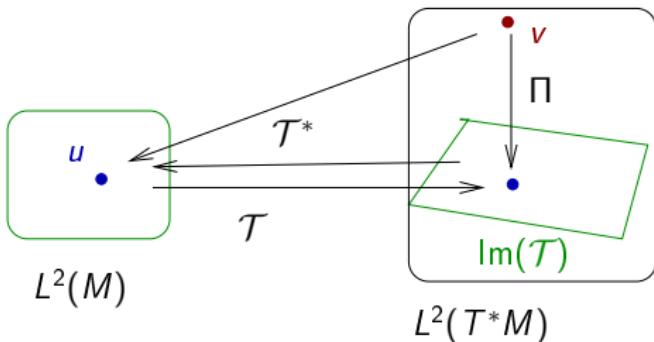
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## Proof of $\mathcal{T}^*\mathcal{T} = \text{Id}_{S(\mathbb{R})}$ for example 1. “Bargman transform”

Let  $\rho = (x, \xi) \in \mathbb{R}^2$ , metric  $g = dx^2 + d\xi^2$  (with  $\alpha = 0$ ), and **wave-packet is the Gaussian function**

$$\varphi_\rho(y) = a \exp(i\xi y) \exp\left(-\frac{1}{2}|y-x|^2\right), \quad a = \pi^{-1/4}. \quad (3)$$

Recall that  $(\mathcal{T}u)(\rho) := \langle \varphi_\rho | u \rangle_{L^2(\mathbb{R})}$  and  $(\mathcal{T}^*v)(y) = \int v(x, \psi) \varphi_{x, \xi}(y) \frac{dx d\xi}{2\pi}$ . Then

$$\begin{aligned} \langle \delta_{y'} | \mathcal{T}^* \mathcal{T} \delta_y \rangle &= \int_{\mathbb{R}^2} \langle \delta_{y'} | \mathcal{T}^* \delta_{x, \xi} \rangle \langle \delta_{x, \xi} | \mathcal{T} \delta_y \rangle dx d\xi \\ &= \int \varphi_{x, \xi}(y') \overline{\varphi_{x, \xi}(y)} \frac{dx d\xi}{2\pi} \\ &= \frac{a^2}{2\pi} \int e^{i\xi(y'-y) - \frac{1}{2}|x-y'|^2 - \frac{1}{2}|x-y|^2} dx d\xi \end{aligned}$$

Since  $\int_{\mathbb{R}} e^{i\xi(y'-y)} d\xi = (2\pi) \delta(y' - y)$  and  $\int_{\mathbb{R}} e^{-|x|^2} dx = \pi^{1/2}$  we get

$$\begin{aligned} \langle \delta_{y'} | \mathcal{T}^* \mathcal{T} \delta_y \rangle &= a^2 \delta(y' - y) \int e^{-|x-y|^2} dx \\ &= a^2 \pi^{1/2} \delta(y' - y) \\ &= \delta(y' - y) \quad \Leftrightarrow \mathcal{T}^* \mathcal{T} = \text{Id}_{S(\mathbb{R})} \end{aligned}$$

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$$\begin{aligned} \langle \delta_{y'} | \mathcal{T}^* \mathcal{T} \delta_y \rangle &= \int_{\mathbb{R}^2} \langle \delta_{y'} | \mathcal{T}^* \delta_{x, \xi} \rangle \langle \delta_{x, \xi} | \mathcal{T} \delta_y \rangle dx d\xi \\ &= \int \varphi_{x, \xi}(y') \overline{\varphi_{x, \xi}(y)} \frac{dx d\xi}{2\pi} \\ &= \frac{a^2}{2\pi} \int e^{i\xi(y'-y) - \frac{1}{2}|x-y'|^2 - \frac{1}{2}|x-y|^2} dx d\xi \end{aligned}$$

Since  $\int_{\mathbb{R}} e^{i\xi(y'-y)} d\xi = (2\pi) \delta(y' - y)$  and  $\int_{\mathbb{R}} e^{-|x|^2} dx = \pi^{1/2}$  we get

$$\begin{aligned} \langle \delta_{y'} | \mathcal{T}^* \mathcal{T} \delta_y \rangle &= a^2 \delta(y' - y) \int e^{-|x-y|^2} dx \\ &= a^2 \pi^{1/2} \delta(y' - y) \\ &= \delta(y' - y) \quad \Leftrightarrow \mathcal{T}^* \mathcal{T} = \text{Id}_{S(\mathbb{R})} \end{aligned}$$

## Proof of $\mathcal{T}^*\mathcal{T} = \text{Id}_{S(\mathbb{R})}$ for example 1. “Bargman transform”

Let  $\rho = (x, \xi) \in \mathbb{R}^2$ , metric  $g = dx^2 + d\xi^2$  (with  $\alpha = 0$ ), and **wave-packet is the Gaussian function**

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# Trace formula, “uncertainty principle”

Let  $A \subset T^*M$ , measurable set and  $\chi_A$  its characteristic function. Let

$$\text{Op}(\chi_A) := \mathcal{T}^* \chi_A \mathcal{T} : L^2(M) \rightarrow L^2(M)$$

called “**Toeplitz quantization of  $\chi_A$** ” that restricts functions to “their components on  $A$ ”.

One has  $\|\text{Op}(\chi_A)\|_{L^2} \leq 1$ ,  $\text{Op}(\chi_A) \geq 0$  and

$$\begin{aligned}\text{Tr}(\text{Op}(\chi_A)) &= \text{Tr}(\mathcal{T} \mathcal{T}^* \chi_A) = \int_{T^*M} \langle \delta_\rho | \mathcal{T} \mathcal{T}^* \delta_\rho \rangle \chi_A(\rho) d\rho \\ &= \int_{T^*M} \underbrace{\langle \varphi_\rho | \varphi_\rho \rangle}_{\approx 1 \text{ if } |\eta| \gg 1} \frac{1}{(2\pi)^{\dim M}} \chi_A(\rho) d\rho \\ &\approx \frac{1}{(2\pi)^{\dim M}} \text{Vol}(A).\end{aligned}$$

**Interpretation:** the number of independent functions  $u \in C^\infty(M)$  that “can live” on  $A \subset T^*M$  is  $\approx \frac{1}{(2\pi)^{\dim M}} \text{Vol}(A)$ .

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# Propagation of singularities

Consider the lifted transfer operator:

$$\begin{array}{ccc} L^2(M) & \xrightarrow{e^{tX}} & L^2(M) \\ \tau \downarrow & & \tau \downarrow \\ L^2(T^*M) & \xrightarrow{\mathcal{T} e^{tX} \mathcal{T}^*} & L^2(T^*M) \end{array}$$

The flow  $\phi^t$  induces the flow  $\tilde{\phi}^t(y, \eta) = (\phi^{-t}(y), (D\phi^t)_y^* \eta)$  on  $T^*M$ .

**Theorem (fundamental 2. "Propagation of singularities", for any vector field  $X$ )**

$\forall t \geq 0, \forall N > 0, \exists C_{N,t} > 0, \forall \rho, \rho' \in T^*M,$

$$|\langle \delta_{\rho'} | \mathcal{T} e^{tX} \mathcal{T}^* \delta_\rho \rangle_{L^2(T^*M)}| = |\langle \varphi_{\rho'} | e^{tX} \varphi_\rho \rangle_{L^2(M)}| \leq C_{N,t} \left\langle \text{dist}_g(\rho', \tilde{\phi}^t(\rho)) \right\rangle^{-N} \quad (4)$$

- i.e.  $e^{tX}$  is a **Fourier Integral Operator (FIO)**. This is the **fundamental property** for microlocal analysis, i.e.  $e^{tX}$  is a quasi orthogonal sum of finite rank operators parametrized by  $T^*M$ , but for a metric  $g$  well chosen.
- In particular  $t = 0$  gives quasi orthogonality of wave packets (2).

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# Proof of propagation of singularities for example 1

Recall  $X = -x\partial_x$  on  $\mathbb{R}$ . Lifted flow is  $\tilde{\phi}^t(x, \xi) = (e^t x, e^{-t} \xi)$ .

Metric  $g = dx^2 + d\xi^2$ .

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Using (3) and Gaussian integrals we get

$$|\langle \delta_{\rho'} | \mathcal{T} e^{tX} \mathcal{T}^* \delta_\rho \rangle| = \left( \frac{2}{1 + e^{-2t}} \right)^{1/2} \exp \left( -\frac{1}{2(1 + e^{-2t})} \left( \left( \frac{x' - e^t x}{e^t} \right)^2 + (\xi' - e^{-t} \xi)^2 \right) \right)$$

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# Objective: understand $\tilde{\phi}^t : T^*M \rightarrow T^*M$

Recall: for  $X$  Anosov vector field,

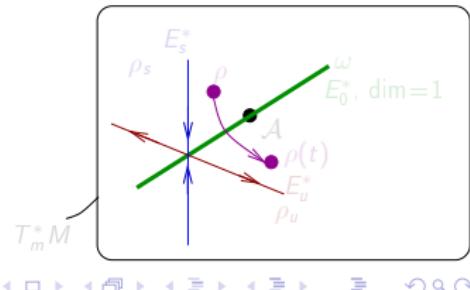
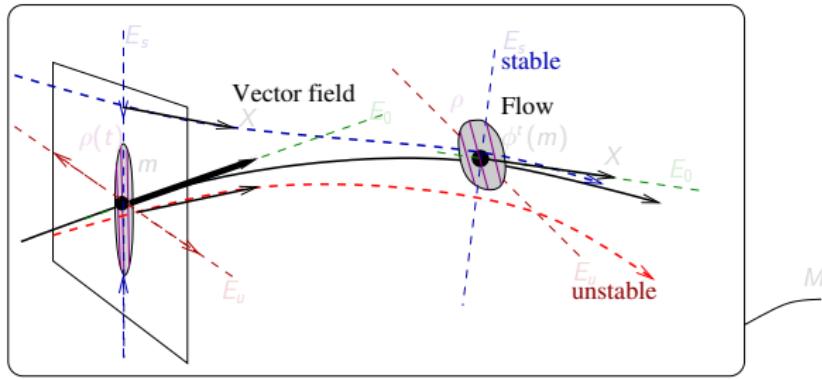
$$TM = E_u \oplus E_s \oplus \underbrace{E_0}_{\mathbb{R}X}$$

Let

$$E_0^* := (E_u \oplus E_s)^\perp = \{\eta \in T^*M, \text{ s.t. } \text{Ker}(\eta) \supset (E_u \oplus E_s)\} = \mathbb{R}\mathcal{A}$$

with Anosov one-form  $\mathcal{A}(X) = 1$ ,  $\text{Ker}\mathcal{A} = E_u \oplus E_s$  that is  $C^{\beta_0}$ .

$$E_s^* := (E_u \oplus E_0)^\perp, \quad E_u^* := (E_s \oplus E_0)^\perp.$$



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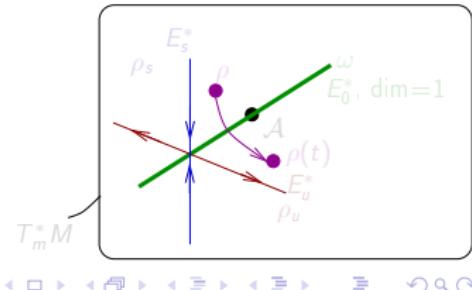
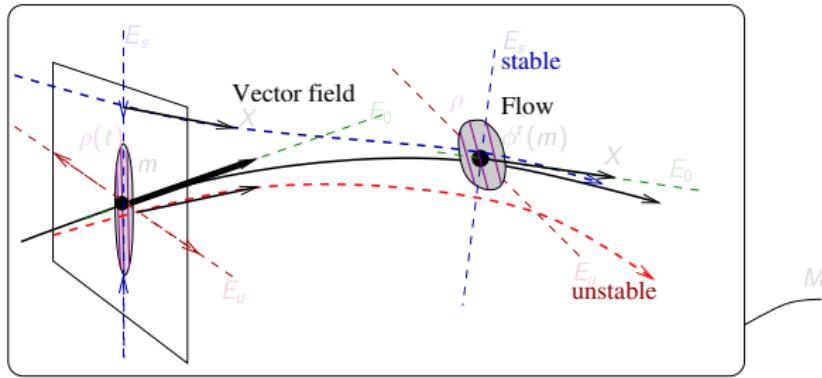
$$TM = E_u \oplus E_s \oplus \underbrace{E_0}_{\mathbb{R}X}$$

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Objective: understand  $\tilde{\phi}^t : T^*M \rightarrow T^*M$

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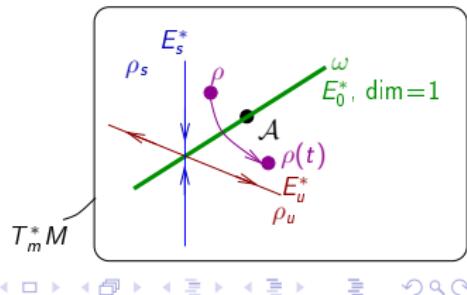
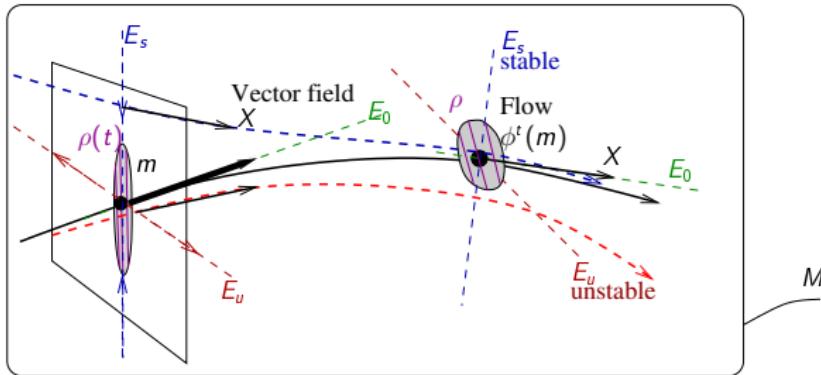
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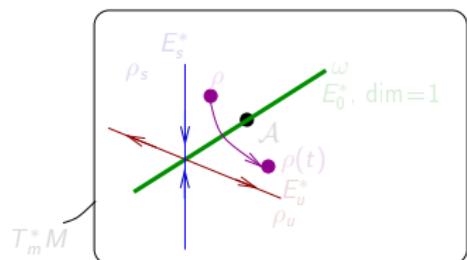
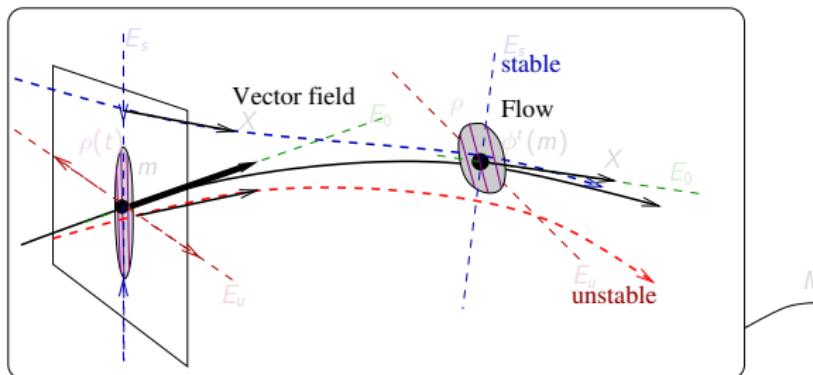
# Scattering in $T^*M$ on the trapped set $E_0^*$

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$$\rho = \rho_u + \rho_s + \omega \mathcal{A}(m) \in T^*M, \quad \omega \in \mathbb{R} \text{ frequency.}$$

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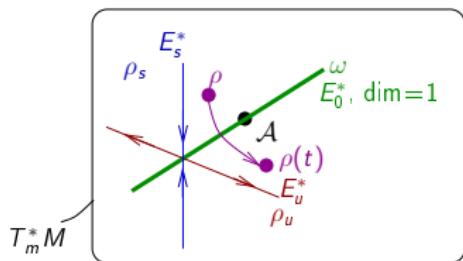
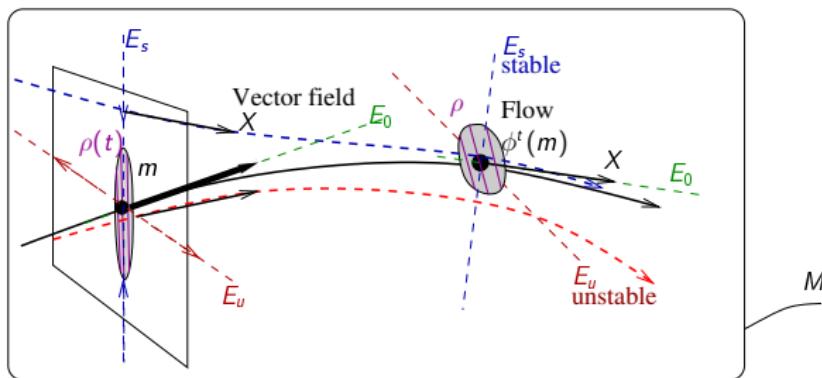
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# Weight (escape) function $W$ in $T^*M$

Let  $\alpha \geq \frac{1}{1+\beta_0}$  so that “the metric absorbs the  $\beta_0$ –Hölder fluctuations of  $E_0^*$ ” (explained later..).

Notation:  $\langle x \rangle := |x|$  if  $|x| \geq 1$ , otherwise  $\langle x \rangle = 1$ .

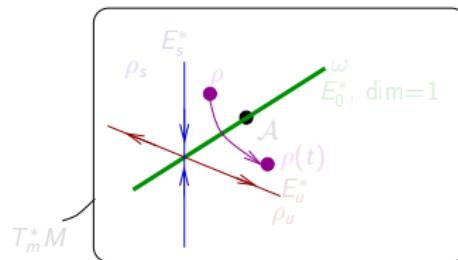
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$$h_\gamma(\rho) = h_0 \left\langle \|\rho_u + \rho_s\|_g \right\rangle^{-\gamma}$$

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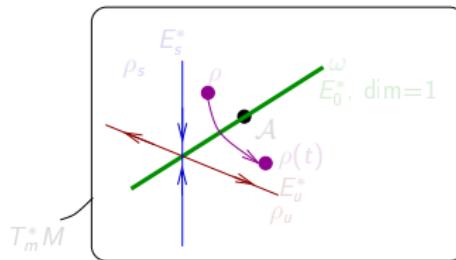
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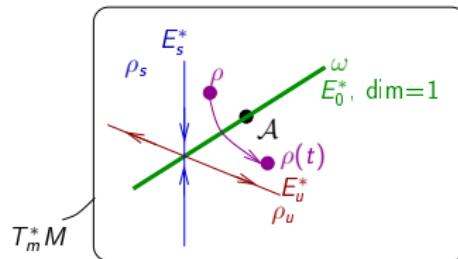
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$$\forall t \geq 0, \quad \frac{1}{C} e^{-(\lambda_{\max} r)t} \leq \frac{W(\tilde{\phi}^t(\rho))}{W(\rho)} \leq \begin{cases} C \\ Ce^{-(\lambda r)t} & \text{if } \text{dist}_g(\rho, E_0^*) > C_t \end{cases} \quad (6)$$

with  $r = R(1 - \gamma)(1 - \alpha)$  (the order of  $W$ ) i.e.

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Recall  $X = -x\partial_x$  on  $\mathbb{R}$ . Lifted flow is  $\tilde{\phi}^t(x, \xi) = (e^t x, e^{-t} \xi)$ .

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# Anisotropic Sobolev space $\mathcal{H}_W(M)$

## Definition

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$$\|u\|_{\mathcal{H}_W(M)} := \|W\mathcal{T}u\|_{L^2(T^*M)},$$

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i.e. we have isometries

$$\mathcal{H}_W(M) \xrightarrow{\mathcal{T}} L^2(T^*M; W^2) \xrightarrow{W} L^2(T^*M)$$

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Good choice of symplectic boxes  $\alpha \geq \frac{1}{1+\beta_0}$  for the phase space metric  $g$

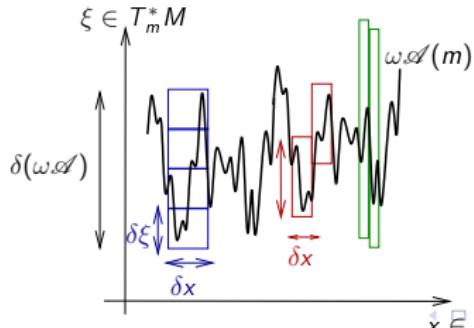
Recall, (for the trapped set  $|\eta| \sim |\omega| \gg 1$ ):

$$g_\rho = \left( \frac{dx}{\langle \omega \rangle^{-\alpha}} \right)^2 + \left( \frac{dz}{\delta} \right)^2 + \left( \frac{d\xi}{\langle \omega \rangle^\alpha} \right)^2 + \left( \frac{d\omega}{\delta^{-1}} \right)^2$$

The trapped set is  $E_0^* = \mathbb{R}_\omega \mathcal{A}$ , with  $\mathcal{A}$  that is  $\beta_0$ -Hölder. Hence variation  $\delta x := \langle \omega \rangle^{-\alpha}$  gives  $\delta(\omega \mathcal{A}) = \omega (\delta x)^{\beta_0} = \omega^{1-\alpha\beta_0}$ .

For **temperate property** of  $W$  one requires

$$\begin{aligned}\delta(\omega\mathcal{A}) \leq \delta\xi &:= \langle\omega\rangle^\alpha \Leftrightarrow \omega^{1-\alpha\beta_0} \leq \omega^\alpha \\ &\Leftrightarrow 1 - \alpha\beta_0 \leq \alpha \\ &\Leftrightarrow \alpha \geq \frac{1}{1 + \beta_0}\end{aligned}$$



# Remark about Hörmander classes of symbols $S_{\rho,\delta}^m$

In transverse direction  $(x, \xi)$ , after smoothing,

$$W \in S_{\rho,\delta}^m \quad \stackrel{\text{def}}{\Leftrightarrow} \quad |\partial_\xi^a \partial_x^b W| \leq C_{a,b} \langle \xi \rangle^{m - \rho|a| + \delta|b|}$$

with

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Ref: **Hörmander's PDE book Vol.3**, or **Lerner's book** p.68, with a conformal “phase’s space metric”

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## PART 3: Some results

Versions of the next Theorem about **discrete RP spectrum** have already been obtained

- For **Anosov diffeomorphisms**, by Blank, Keller, Liverani 2002, Gouëzel, Liverani 2005, Baladi, Tsujii 2005,2008, F-Roy-Sjöstrand 2008.
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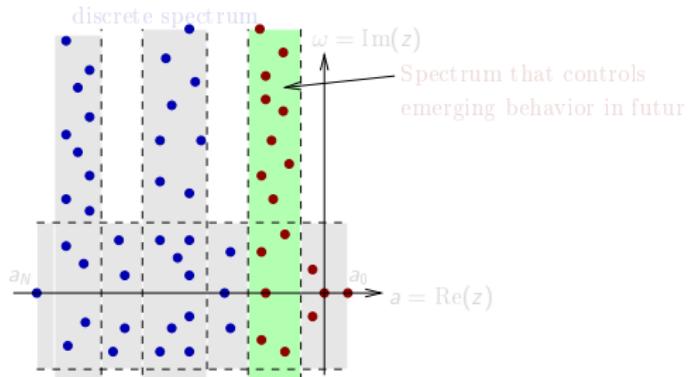
# Question: discrete spectrum of the generator $X_F$ ?

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$\sim \underset{t \rightarrow +\infty}{e^{ta_0} e^{it\omega_0} \Pi_0 + \dots}$  : if  $a_0 > a_{j \neq 0}$ , : "futur emerging behavior"

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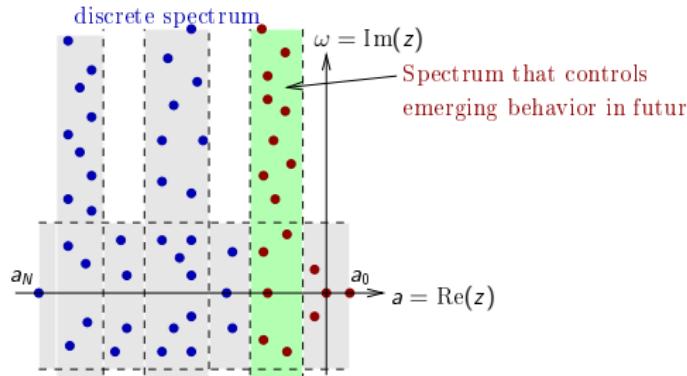
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## Meaning of this internal spectrum? Illustration

- In first approximation, the sea is quite, deep, flat, gentle  $\equiv$  **Equilibrium state  $z_0, z_N$ , dominant.**



- But the behavior at the surface may be furious  $\equiv$  **internal spectrum  $(z_j)_j$**



## Notations

Consider **Anosov vector field**  $X$  on  $M$  and **potential**  $V \in C^\infty(M)$ ,

$$X_F := X + V$$

(or more general  $X_F : C^\infty(M; F) \circlearrowleft$ ).

$$C_{X,V} := \overline{\max} \left( \frac{1}{2} \operatorname{div} X + \operatorname{Re}(V) \right) \quad (9)$$

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## Ruelle Pollicott discrete spectrum of resonances

Recall weight function  $W(\rho) = \frac{\langle h_\gamma(\rho) \|\rho_s\|_g \rangle^R}{\langle h_\gamma(\rho) \|\rho_u\|_g \rangle^R}$  has order  $r = R(1-\gamma)(1-\alpha)$ , i.e.  
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### Theorem (F.-Tsujii 17)

$X_F := X + V$ . The family of operators

$$e^{tX_F} : \mathcal{H}_W(M) \rightarrow \mathcal{H}_W(M), t \in \mathbb{R}$$

form a **strongly continuous group**: for  $t \geq 0$ ,

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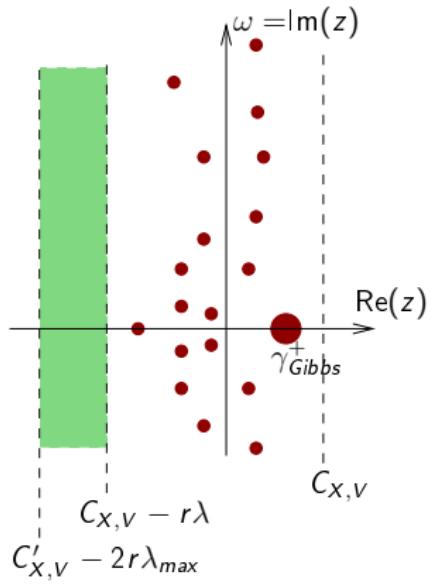
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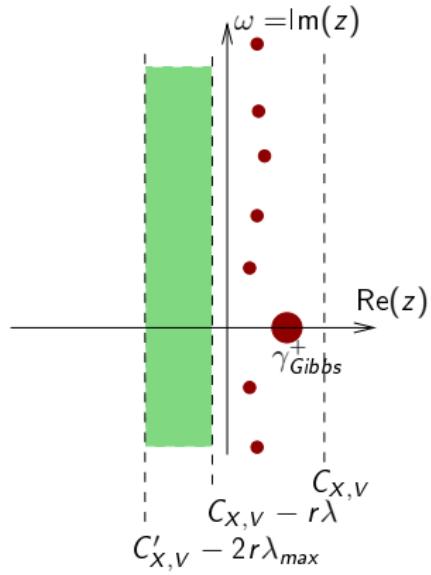
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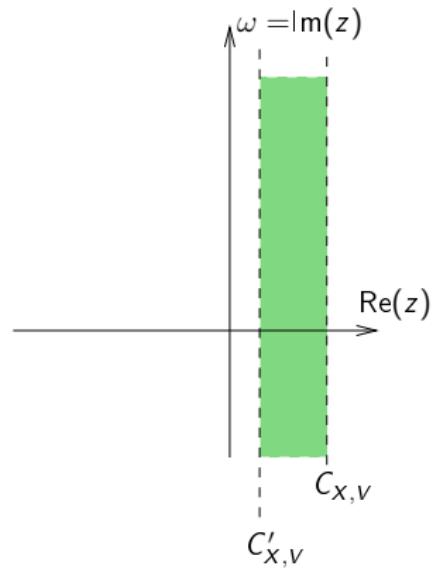
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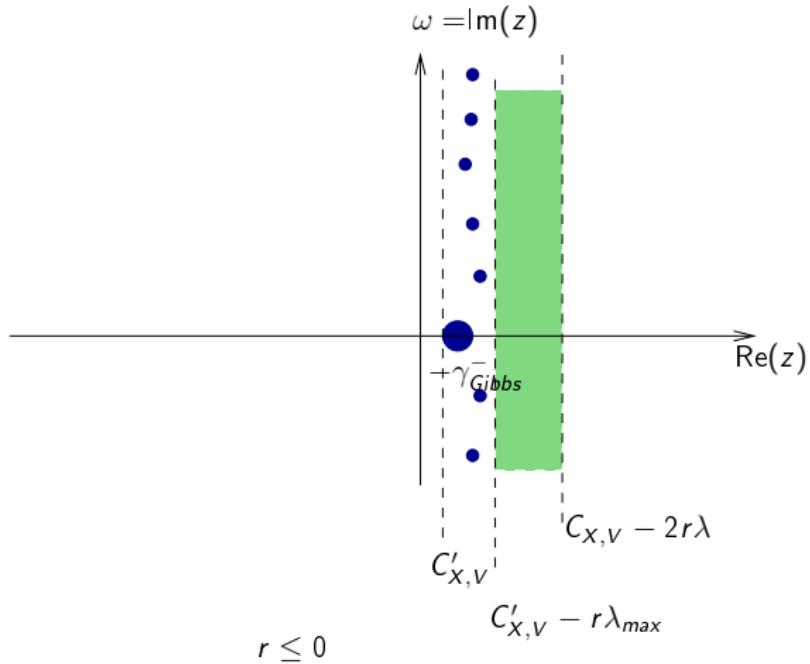
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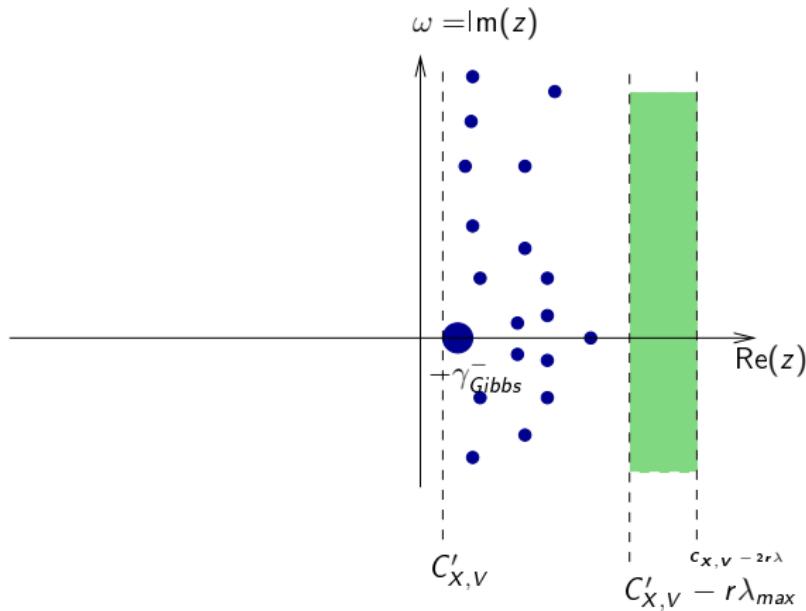


$$r = 0$$

Spectrum of  $X_F$  in  $\mathcal{H}_W(M) = L^2(M)$



Ruelle-Pollicott spectrum  $\sigma_-(X_F)$  of the past dynamics.



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## Remark on past/future spectrum

- If  $V' + \overline{V} + \operatorname{div}X = 0$  then the **future** spectrum  $\lambda_j^+$  of  $A = -X + V$  is related to the **past** spectrum  $\lambda'_j^-$  of  $A' = -X + V'$  by

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# Fractal Weyl law

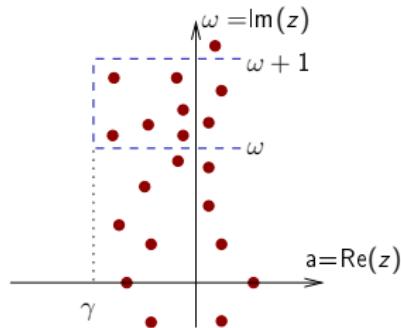
(after J. Sjöstrand 90 for quantum resonances)

Theorem (F.-Tsujii 17. "Upper bound for the density of eigenvalues")

$\forall \gamma \in \mathbb{R}, \exists C > 0, \forall \omega \geq 1,$

$$\#\{z \in \sigma(X); \quad \operatorname{Re}(z) > \gamma, \operatorname{Im}(z) \in [\omega, \omega + 1]\} \leq C |\omega|^{\frac{\dim M - 1}{1 + \beta_0}}.$$

with  $\beta_0 \in ]0, 1]$  is Hölder exponent of  $E_u \oplus E_s$ .



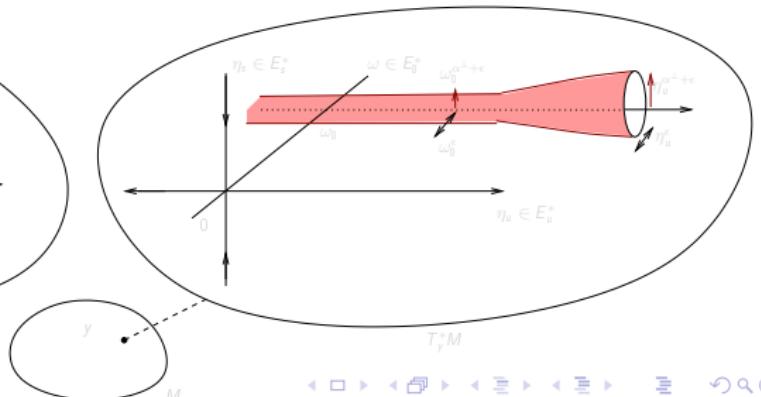
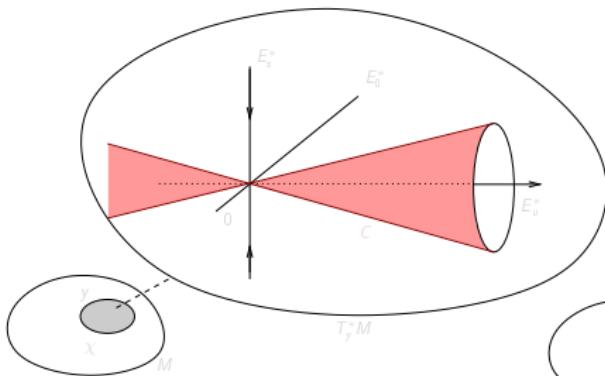
# About the wave front set of resonances

Theorem (F.-Tsujii 17. "Parabolic wave front set")

$\forall C, N, \epsilon, \exists C_N$ , for any (generalized) Ruelle Pollicott eigenfunction  $u$  with  $\operatorname{Re}(z) > -C$  then  $\forall (y, \eta) \in T^* M$ ,

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We choose  $\alpha = \frac{1}{1+\min(\beta_u, \beta_s)}$  (but expect  $\alpha = \frac{1}{1+\beta_u}$ ) so that uncertainty principle absorbs Hölder fluctuations.



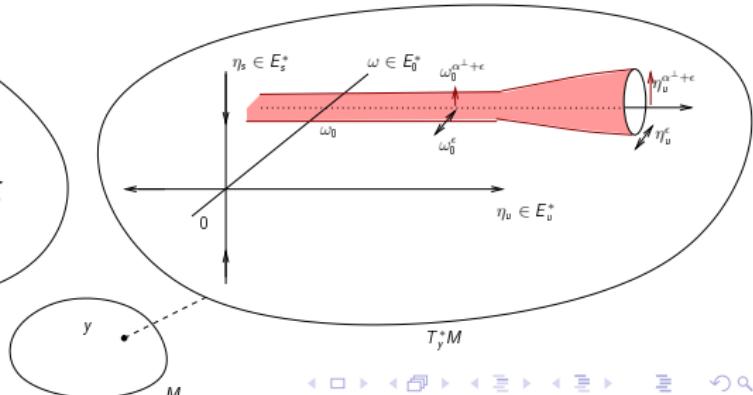
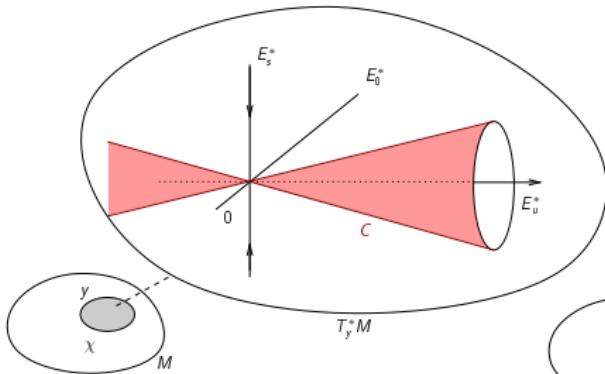
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# Sketch of proof for group property (1)

Recall the isometries:

$$\mathcal{H}_W(M) \xrightarrow{\mathcal{T}} L^2(T^*M; W^2) \xrightarrow{W} L^2(T^*M)$$

The strategy of microlocal analysis is to analyse the conjugated operator  $W\mathcal{T}e^{tX}\mathcal{T}^*W^{-1}$  in  $L^2(T^*M)$  instead of  $e^{tX}$ :

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## Sketch of proof for group property (2)

One has

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for any  $N > 0$ , i.e. decays outside the graph of  $\tilde{\phi}^t$ .

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- By using that  $h_\gamma(\rho) \ll 1$  in  $W(\rho)$  one can **improve the above estimate** and get the bound:  $\|e^{tX}\|_{H_W(M)} \leq C e^{(C_{X,V} + o(1))t}$ . Similarly for  $t \leq 0$ .

## Sketch of proof for “discrete RP spectrum” and Weyl law

Recall decay property (6):  $\frac{W(\tilde{\phi}^t(\rho))}{W(\rho)} \leq C e^{-(\lambda r)t}$  if  $\text{dist}_g(\rho, E_0^*) \geq \sigma \gg 1$ . Let  $\chi_\sigma(\rho)$  characteristic function for  $\text{dist}_g(\rho, E_0^*) \geq \sigma$ .

If we repeat the previous argument, the norm decays very fast far from the trapped set:

$$\|W \mathcal{T} e^{tX} \mathcal{T}^* W^{-1} \chi_\sigma\|_{L^2(T^* M)} \leq C e^{(C_{X,V} - \lambda r + o(1))t}.$$

Near the trapped set  $E_0^*$ , let

$$A := \{\rho \in T^* M, \quad \text{dist}_g(\rho, E_0^*) \leq \sigma, \quad \omega(\rho) \in [\omega, \omega + 1]\}.$$

Using **trace formula** (and Jensen inequality..), the **density of RP spectrum** is then

$$\begin{aligned} \#\{z \in \sigma(X), \quad \text{Im}(z) \in [\omega, \omega + 1], \text{Re}(z) \geq -\gamma\} &\lesssim \text{Tr}(\mathcal{T}^* \chi_A \mathcal{T}) \\ &\asymp \text{Vol}(A) \\ &\asymp (\sigma \omega^\alpha)^{\dim M - 1} \\ &\asymp \omega^{\frac{\dim M - 1}{1 + \beta_0}} \end{aligned}$$

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Thank you for your attention

