# The Ising Model's Low Temperature Representation and Phase Transition

## Spontaneous symmetry breaking at low temperatures

It is shown that for the nearest-neighbour Ising model on  $\mathbb{Z}^d$ , the critical inverse temperature  $\beta_c(d)$  satisfies  $\beta_c(d) < \infty$ , for all dimensions  $d \geq 2$ , implying that the model undergoes a phase transition at low temperatures.

### The Ising model

For a finite subset  $\Lambda \subset \mathbb{Z}^d$ , the set of *configura*tions is  $\Omega_{\Lambda} := \{-1, 1\}^{\Lambda}$ . The spin  $\sigma_i$  at vertex  $i \in \mathbb{Z}^d$  is the random variable giving the value of the configuration at vertex i.

The interactions among spins in the nearestneighbour model are only between neighbouring spins and favour agreement of spin values.

The distribution of configurations, with (+)boundary, is given by the Gibbs measure  $\mu_{\Lambda;\beta}^+$ . One may extend  $\mu_{\Lambda:\beta}^+$  on all of  $\mathbb{Z}^d$  to the *infinite* volume Gibbs state, under which, expectation shall be denoted  $\langle \cdot \rangle_{\beta}^+$ .

### Phase transition and the critical temperature

We say that a *phase transition* occurs at temperature  $\beta$  if at least two distinct Gibbs states can be constructed at  $\beta$ . The spontaneous magnetisation  $m^*(\beta)$  is defined to be  $m^*(\beta) := \langle \sigma_0 \rangle_{\beta}^+$ and the critical inverse temperature  $\beta_c(d) :=$  $\inf\{\beta \ge 0 \mid m^*(\beta) > 0\}.$ 

#### Low-temperature representation

A low temperature favors the alignment of nearest-neighbour spins. Therefore the contours, which are lines that separate regions of + and - spins, should be sparse. This geometric observation is used, for the case d=2, to derive an expression of  $\mu_{\Lambda:\beta}^+$  called the low-temperature representation.

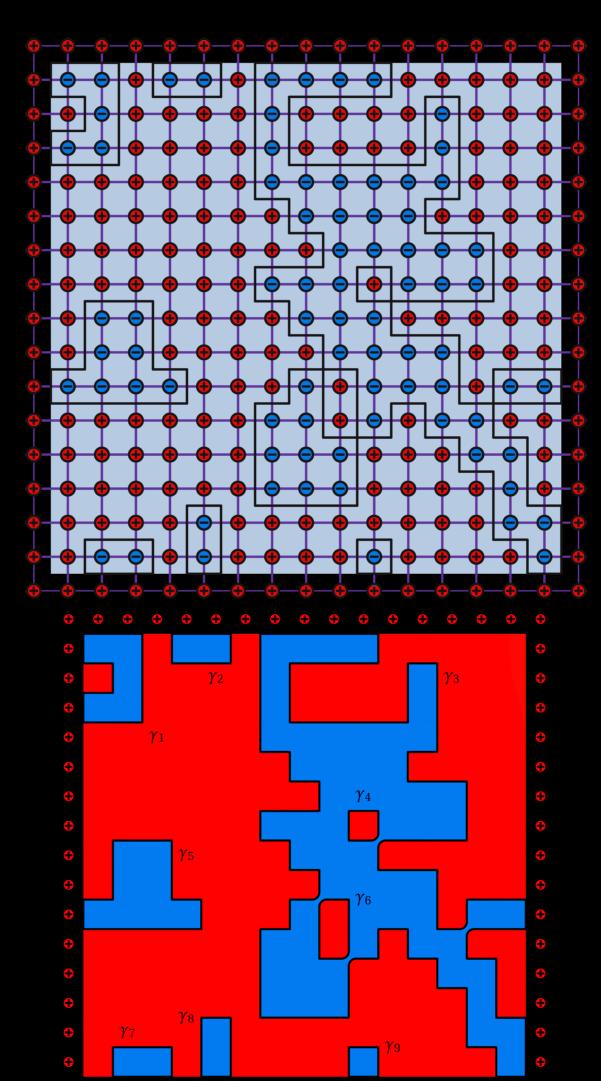


Fig. 1: Given the contrours of a configuration, which are paths on the dual lattice, it may be uniquely reconstructed.

### Peierls' argument

Let  $B_n = \{-n, ..., n\}^2$ . For any  $\omega \in \Omega_{B_n}^+$  let  $\Gamma(\omega)$  be the set of contours of  $\omega$  as in Fig.1. For all  $\beta > 0$  and any contour  $\gamma$ ,

$$\mu_{B_n;\beta}^+(\Gamma \ni \gamma) \le e^{-2\beta|\gamma|}.$$

#### Extension to higher dimensions

Using the two-dimensional low temperature representation and Peierls' argument, it is shown that  $\beta_c(2) < \infty$ . This analysis is extended for all  $d \geq 3$  by embedding  $\mathbb{Z}^d$  into  $\mathbb{Z}^{d+1}$  and using the GKS inequalities to show that  $\beta_c(d)$  is nonincreasing in d.

#### References

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